

NONLINEAR WAVES AND SOLITONS

Antonio Degasperis

Physics Department
University of Rome *La Sapienza*

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CONTENTS

- **INTRODUCTION: linear and nonlinear waves**
- Multiscale perturbation method
- Integrable wave equations and spectral theory
- Conservation laws
- Darboux transformations and soliton solutions

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INTRODUCTION: LINEAR WAVES

NOTATION: t = time , x = space , $u_t = \partial u / \partial t$, $u_x = \partial u / \partial x$, etc.

LINEAR wave equations in 1+1 dimensions:

TRANSLATION WAVES

$$u_t + v u_x = 0 , \quad u = u(x, t) = u_0(x - vt)$$

DISPERSIVE WAVES

$$u_t + i\omega(-i\frac{\partial}{\partial_x})u = 0 , \quad u(x, 0) = u_0(x)$$

$$u(x, t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} dk \hat{u}_0(k) e^{i[kx - \omega(k)t]}$$

$$u_0(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} dk \hat{u}_0(k) e^{ikx}$$

$\omega = \omega(k)$ = dispersion relation

INTRODUCTION: EXAMPLES OF LINEAR WAVE EQUATIONS

- Schrödinger equation: $\omega(k) = k^2$

$$iu_t + u_{xx} = 0$$

- Linearized Korteweg-de Vries equation: $\omega(k) = k^3$

$$u_t - u_{xxx} = 0$$

- Linearized Benjamin-Ono equation: $\omega(k) = \text{sign}(k)k^2$

$$u_t - Hu_{xx} = 0 , \quad Hf(x) = \frac{1}{\pi} P \int_{-\infty}^{+\infty} dy \frac{f(y)}{y-x}$$

INTRODUCTION: CONSERVATION LAWS

LINEAR WAVE EQUATIONS HAVE INFINITELY MANY CONSERVATION LAWS

LOCAL :

$$\rho_t(x, t) + j_x(x, t) = 0$$

GLOBAL :

$$C = \int_{-\infty}^{+\infty} dx \rho(x, t) , \quad \frac{dC}{dt} = 0$$

EXERCISE: FIND 3 OR 4 LOCAL CONSERVATION LAWS OF THE WAVE EQUATIONS DISPLAYED ABOVE

INTRODUCTION: NONLINEAR WAVE EQUATIONS

$$Lu = N(u)$$

GENERICALLY A NONLINEAR WAVE EQUATION POSSESSES ONLY FEW CONSERVATION LAWS

EXCEPTIONALLY A WAVE EQUATION HAS INFINETELY MANY CONSERVATION LAWS AND IT IS INTEGRABLE

EXAMPLES:

$$u_t + u_{xxx} = P(u)u_x , \quad P(u) = \sum_{n=0}^N c_n u^n$$

THEOREM: THIS EQUATION IS INTEGRABLE IF AND ONLY IF

$$N \leq 2$$

$$P(u) = c_0 + c_1 u + c_2 u^2$$

INTRODUCTION: NONLINEAR WAVE EQUATIONS CONT.

C-INTEGRABILITY: LINEARISATION VIA A CHANGE OF VARIABLES

ECKHAUS EQUATION: $iu_t + u_{xx} = [2(|u|^2)_x + |u|^4] u$

$$u(x, t) \rightarrow v(x, t) = u(x, t) \exp\left[\int_{x_0}^x dy |u(y, t)|^2\right]$$
$$iv_t + v_{xx} = 0$$

EXERCISES:

- ① FIND THE INVERSE TRANSFORMATION $v(x, t) \rightarrow u(x, t)$
- ② FIND EXPLICIT SOLUTIONS OF THE ECKHAUS EQUATION
- ③ DISCUSS THE STABILITY OF THE PLANE WAVE SOLUTION

INTRODUCTION: NONLINEAR WAVE EQUATIONS CONT.

S-INTEGRABILITY: SPECTRAL ANALYSIS (NONLINEAR FOURIER TRANSFORM)

KORTEWEG- DE VRIES EQUATION (KdV): $u_t + u_{xxx} = c_0 u_x + c_1 u u_x$

MODIFIED KORTEWEG- DE VRIES EQUATION (mKdV):

$$u_t + u_{xxx} = c_0 u_x + c_2 u^2 u_x$$

NONLINEAR SCHROEDINGER EQUATION (NLS):

$$iu_t + u_{xx} = 2\eta|u|^2 u$$

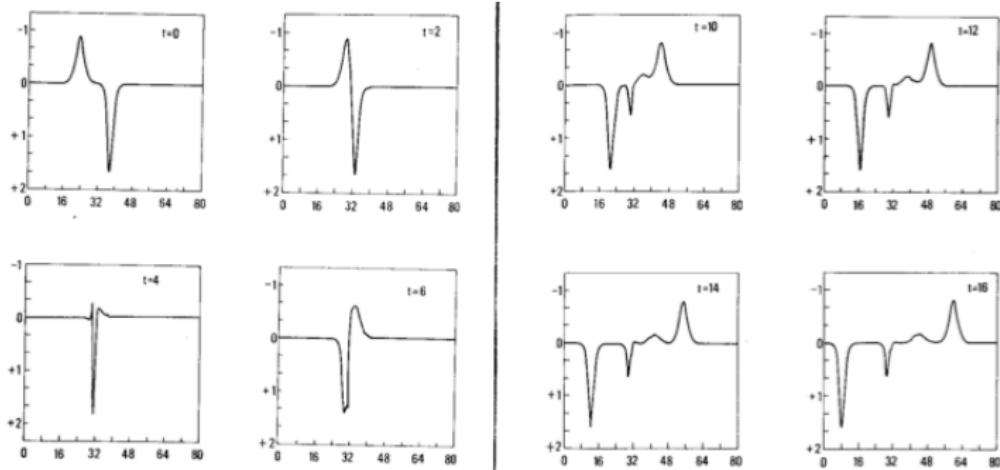
$$\Psi_x = \begin{pmatrix} -ik & \eta u^* \\ u & ik \end{pmatrix} \Psi , \quad \Psi = \begin{pmatrix} \psi_1(x, t, k) \\ \psi_2(x, t, k) \end{pmatrix} , \quad u = u(x, t)$$

$$\Psi_t = \begin{pmatrix} 2ik^2 + i\eta|u|^2 & -2k\eta u^* - i\eta u_x^* \\ -2ku + iu_x & -2ik^2 - i\eta|u|^2 \end{pmatrix} \Psi$$

$$\Psi_{xt} = \Psi_{tx}$$

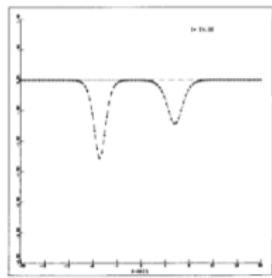
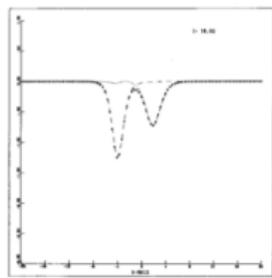
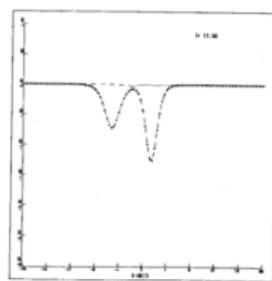
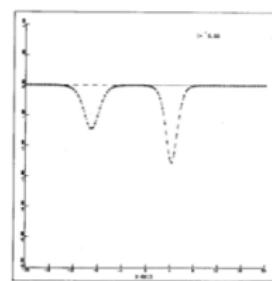
INTRODUCTION: NONLINEAR WAVE EQUATIONS CONT.

REGULARIZED LONG WAVE EQUATION $u_t + u_x - u_{xxt} = 6uu_x$



INTRODUCTION: NONLINEAR WAVE EQUATIONS CONT.

KDV EQUATION: $u_t + u_{xxx} = 6uu_x$



INTRODUCTION: NONLINEAR WAVE EQUATIONS CONT.

HISTORICAL NOTES

1845.....SCOTT-RUSSELL

1895.....KORTEWEG-DE VRIES

1955.....FERMI-PASTA-ULAM

1965.....ZABUSKY-KRUSKAL

1967.....GARDNER-GREENE-KRUSKAL-MIURA

1971.....ZAKHAROV-SHABAT

MULTISCALE ODE 1

POINCARÉ-LINDSTEDT

$$\ddot{q} + \omega_0^2 q = c_2 q^2 + c_3 q^3 + \dots , \quad q = q(t, \epsilon)$$

$$q(0, \epsilon) = \epsilon , \quad \dot{q}(0, \epsilon) = 0$$

$$q(t, \epsilon) = q\left(t + \frac{2\pi}{\omega(\epsilon)}, \epsilon\right) , \quad \theta = \omega(\epsilon)t , \quad q(t, \epsilon) = f(\theta, \epsilon)$$

$$\omega^2(\epsilon)f'' + \omega_0^2 f = c_2 f^2 + c_3 f^3 + \dots , \quad f(0, \epsilon) = \epsilon , \quad f'(0, \epsilon) = 0$$

$$\omega(\epsilon) = \omega_0 + \omega_1 \epsilon + \omega_2 \epsilon^2 + \dots , \quad f(\theta, \epsilon) = \epsilon f_1(\theta) + \epsilon^2 f_2(\theta) + \dots$$

MULTISCALE ODE 1

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MULTISCALE ODE 1

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$$\omega^2(\epsilon)f'' + \omega_0^2 f = c_2 f^2 + c_3 f^3 + \dots , \quad f(0, \epsilon) = \epsilon , \quad f'(0, \epsilon) = 0$$

$$\omega(\epsilon) = \omega_0 + \omega_1 \epsilon + \omega_2 \epsilon^2 + \dots , \quad f(\theta, \epsilon) = \epsilon f_1(\theta) + \epsilon^2 f_2(\theta) + \dots$$

MULTISCALE ODE 2

$$O(\epsilon) : \quad f_1'' + f_1 = 0 , \quad f_1(0) = 1 , \quad f_1'(0) = 0$$

$$O(\epsilon^n) : \quad \begin{cases} f_n'' + f_n = \{-n, -n+1, \dots, -1, 0, 1, \dots, n-1, n\} \\ f_n(0) = 0 , \quad f_n'(0) = 0 \end{cases}$$

$$\{-n, -n+1, \dots, -1, 0, 1, \dots, n-1, n\} = \sum_{\alpha=-n}^n C_n^\alpha e^{i\alpha\theta}$$

SECULARITY EFFECT:

$$\phi'' + \phi = Ce^{i\theta} , \quad \phi(\theta) = Ae^{i\theta} + Be^{-i\theta} + \frac{1}{2i}C\theta e^{i\theta}$$

SECULARITY-FREE CONDITION: $C_n^{\pm 1} = 0$

FREQUENCY RENORMALIZATION EXERCISE: COMPUTE $\omega(\epsilon)$ UP TO $O(\epsilon^2)$ ANSWER: $\omega_1 = 0 , \quad \omega_2 = -(10c_2^2 + 9\omega_0^2 c_3)/24\omega_0^3$

MULTISCALE ODE 3

$$\theta = \omega(\epsilon)t = \omega_0 t + \omega_1 \epsilon t + \omega_2 \epsilon^2 t + \dots \quad , \quad t_n = \epsilon^n t$$

$$\theta = \omega_0 t + \omega_1 t_1 + \omega_2 t_2 + \dots$$

$$q(t, \epsilon) = \sum_{n=1}^{\infty} \sum_{\alpha=-n}^n \epsilon^n \exp(i\alpha\theta) f_n^{(\alpha)}$$

$$q(t, \epsilon) = \sum_{n=1}^{\infty} \sum_{\alpha=-n}^n \epsilon^n E^\alpha q_n^{(\alpha)}(t_1, t_2, \dots) \quad , \quad E \equiv \exp(i\omega_0 t)$$

$$d(E^\alpha q_n^{(\alpha)})/dt = E^\alpha (i\alpha\omega_0 + \epsilon \partial/\partial t_1 + \epsilon^2 \partial/\partial t_2 + \dots) q_n^{(\alpha)}$$

INGREDIENTS :

- EXPANSION IN POWERS OF ϵ
- EXPANSION IN POWERS OF $\exp(i\omega_0 t)$
- INTRODUCTION OF SLOW TIMES t_1, t_2, t_3, \dots

$$\begin{cases} Du = F[u, u_x, u_{xx}, \dots] , \quad u = u(x, t) \\ D = \partial/\partial t + i\omega(-i\partial/\partial x) \\ \omega(k) = \sum_{m=0} a_{2m+1} k^{2m+1} \end{cases}$$

ONE QUASI-MONOCHROMATIC PLANE WAVE

$$u(x, t) \simeq \Delta k \int_{-\infty}^{+\infty} d\eta A(\eta) \exp\{i[x(k+\eta\Delta k)-t\omega(k+\eta\Delta k)]\} + c.c.$$

$$\omega(k + \epsilon\eta k) = \sum_{n=0}^{\infty} \omega_n \eta^n k^n \epsilon^n , \quad \epsilon \equiv \Delta k/k$$

$$u(x, t) \simeq \epsilon E(x, t) u^{(1)}(\xi, t_1, t_2, \dots) + c.c. , \quad E(x, t) \equiv \exp[i(kx - \omega t)],$$
$$\xi \equiv \epsilon x , \quad t_n \equiv \epsilon^n t$$

HARMONIC EXPANSION

$$u(x, t) = \sum_{\alpha=-\infty}^{+\infty} u^{(\alpha)}(\xi, t_1, t_2, \dots) E^\alpha(x, t) , \quad u^{(\alpha)*} = u^{(-\alpha)}$$

$$\partial_x \rightarrow \partial_x + \epsilon \partial_\xi , \quad \partial_t \rightarrow \partial_t + \epsilon \partial_{t_1} + \epsilon^2 \partial_{t_2} + \dots$$

$$D[u^{(\alpha)} E^\alpha] = E^\alpha D^{(\alpha)} u^{(\alpha)}$$

$$F[u, u_x, u_{xx}, \dots] = \sum_{\alpha=-\infty}^{+\infty} F^{(\alpha)}[u^{(\beta)}, u_\xi^{(\beta)}, u_{\xi\xi}^{(\beta)}, \dots] E^\alpha$$

$$D^{(\alpha)} u^{(\alpha)} = F^{(\alpha)}$$

EPSILON EXPANSION

$$D^{(\alpha)} = -i\alpha\omega + \epsilon\partial_{t_1} + \epsilon^2\partial_{t_2} + \dots + i\omega(k - i\epsilon\partial_\xi)$$

$$D^{(\alpha)} = D_0^{(\alpha)} + \epsilon D_1^{(\alpha)} + \epsilon^2 D_2^{(\alpha)} + \dots$$

$$D_0^{(\alpha)} = i[\omega(\alpha k) - \alpha\omega(k)], \quad D_n^{(\alpha)} = \partial_{t_n} - (-i)^{n+1}\omega_n(\alpha k)\partial_\xi^n, \quad n \geq 1$$

$$u^{(\alpha)} = \sum_{n=1} \epsilon^n u^{(\alpha)}(n), \quad |\alpha| > 1$$

assumption : $F \rightarrow -F$ if $u \rightarrow -u \implies u^{(2\alpha)} = 0$

$$F^{(\alpha)} = \epsilon^3 F_3^{(\alpha)} + \epsilon^4 F_4^{(\alpha)} + \dots$$

$$D_0^{(\alpha)} u^{(\alpha)}(n+1) + D_1^{(\alpha)} u^{(\alpha)}(n) + D_2^{(\alpha)} u^{(\alpha)}(n-1) + \dots + D_n^{(\alpha)} u^{(\alpha)}(1) = F_{n+1}^{(\alpha)}$$

MULTISCALE PDE 4

SLAVE HARMONICS : $D_0^{(\alpha)} = i[\omega(\alpha k) - \alpha\omega(k)] \neq 0$

NO RESONANCE CONDITION : $[\omega(\alpha k) - \alpha\omega(k)] = 0$ IFF $\alpha = \pm 1$

FUNDAMENTAL HARMONIC : $\alpha = 1$

$$O(\epsilon^2) : (\partial_{t_1} + \omega_1 \partial_\xi) u^{(1)}(1) = 0, \quad u^{(1)}(1) = u^{(1)}(1)(\xi - \omega_1 t_1, t_2, t_3, \dots)$$

$$O(\epsilon^3) : (\partial_{t_1} + \omega_1 \partial_\xi) u^{(1)}(2) = -[(\partial_{t_2} - i\omega_2 \partial_\xi^2) u^{(1)}(1) - F_3^{(1)}]$$

SECULARITY OR RESONANCE PHENOMENON

$$(\partial_t - M)v(t) = w(t), \quad (\partial_t - M)w(t) = 0 \implies v(t) = v_0(t) + tw(t)$$

SECULARITY FREE CONDITION : $w(t) = 0$

$$(\partial_{t_2} - i\omega_2 \partial_\xi^2) u^{(1)}(1) - F_3^{(1)} = 0$$

NLS EQUATION : $i\partial_{t_2} u^{(1)}(1) + \omega_2 \partial_\xi^2 u^{(1)}(1) = 2c|u^{(1)}(1)|^2 u^{(1)}(1)$

ONE RESONANCE : SECOND HARMONIC GENERATION

$$\omega(2k) - 2\omega(k) = 0$$

assumption : $F^{(\alpha)} = \epsilon^2 F_2^{(\alpha)} + \epsilon^3 F_3^{(\alpha)} + \epsilon^4 F_4^{(\alpha)} + \dots$

technical assumptions (check it!) : $F_2^{(\alpha)} = \partial_x \hat{F}_2^{(\alpha)}$, $\omega(0) = 0$

$$D_1^{(1)} u_1^{(1)} = c_1 u_1^{(2)} u_1^{(-1)}, \quad D_1^{(2)} u_1^{(2)} = c_2 u_1^{(1)2}$$

$$\text{SHG} \begin{cases} \partial_t E_1 + v_1 \partial_x E_1 = c_1 E_2 E_1^*, & v_1 = \omega_1(k_1) \\ \partial_t E_2 + v_2 \partial_x E_2 = c_2 E_1^2 & , v_2 = \omega_1(k_2) \end{cases}$$

RED \implies VIOLET

TWO QUASI-MONOCHROMATIC PLANE WAVES

$$u(x, t) =$$

$$\sum_{\alpha_1=-\infty}^{+\infty} \sum_{\alpha_2=-\infty}^{+\infty} u^{(\alpha_1, \alpha_2)}(\xi, t_1, t_2, \dots) \exp\{i[(\alpha_1 k_1 + \alpha_2 k_2)x - (\alpha_1 \omega_1 + \alpha_2 \omega_2)t]\}$$

$$\omega_1 = \omega(k_1), \quad \omega_2 = \omega(k_2)$$

NO RESONANCE :

$$\omega(\alpha_1 k_1 + \alpha_2 k_2) = \alpha_1 \omega(k_1) + \alpha_2 \omega(k_2)$$

IFF $\alpha_1 = \pm 1, \alpha_2 = 0$ OR $\alpha_1 = 0, \alpha_2 = \pm 1$

$$\omega_1(k_1) \neq \omega_1(k_2)$$

$$\text{NLS EQ. : } i\partial_{t_2} u^{(1,0)}(1) + \omega_2(k_1) \partial_\xi^2 u^{(1,0)}(1) = 2c_{10} |u^{(1,0)}(1)|^2 u^{(1,0)}(1)$$

$$\text{NLS EQ. : } i\partial_{t_2} u^{(0,1)}(1) + \omega_2(k_2) \partial_\xi^2 u^{(0,1)}(1) = 2c_{01} |u^{(0,1)}(1)|^2 u^{(0,1)}(1)$$

STRONG RESONANCE : $\omega_1(k_1) = \omega_1(k_2)$

$$\text{VNLS} \left\{ \begin{array}{l} i\partial_{t_2} u^{(1,0)}(1) + \omega_2(k_1) \partial_\xi^2 u^{(1,0)}(1) = 2(c_{10}|u^{(1,0)}(1)|^2 \\ \quad + \gamma_{10}|u^{(0,1)}(1)|^2)u^{(1,0)}(1) \\ i\partial_{t_2} u^{(0,1)}(1) + \omega_2(k_2) \partial_\xi^2 u^{(0,1)}(1) = 2(c_{01}|u^{(0,1)}(1)|^2 \\ \quad + \gamma_{01}|u^{(1,0)}(1)|^2)u^{(0,1)}(1) \end{array} \right.$$

MULTISCALE PDE 8

ONE RESONANCE AT $\alpha_1 = 1$, $\alpha_2 = 1$: $\omega(k_1 + k_2) = \omega(k_1) + \omega(k_2)$

THREE WAVE RESONANT INTERACTION

$$\begin{cases} k_3 = k_1 + k_2 \\ \omega_3 = \omega_1 + \omega_2 \end{cases}$$

technical assumptions (check it!) : $F_2 = \partial_x \hat{F}_2$, $\omega(0) = 0$

at $O(\epsilon^2)$
$$\begin{cases} D_1^{(1,0)} u_1^{(1,0)} = c_1 u_1^{(1,1)} u_1^{(0,-1)} \\ D_1^{(0,1)} u_1^{(0,1)} = c_2 u_1^{(1,1)} u_1^{(-1,0)} \\ D_1^{(1,1)} u_1^{(1,1)} = c_3 u_1^{(1,0)} u_1^{(0,1)} \end{cases}$$

3WRI
$$\begin{cases} \partial_t E_1 + v_1 \partial_x E_1 = c_1 E_3 E_2^*, & v_1 = \omega_1(k_1) \\ \partial_t E_2 + v_2 \partial_x E_2 = c_2 E_3 E_1^*, & v_2 = \omega_1(k_2) \\ \partial_t E_3 + v_3 \partial_x E_3 = c_3 E_1 E_2, & v_3 = \omega_1(k_3) \end{cases}$$

EXERCISES :

- ① DERIVE THE NLS EQUATION FROM THE PDE

$$u_t + au_{xxx} + bu_x = f_1 uu_x + f_2 u^2 u_x$$

BY MULTISCALE METHOD AND TELL IF IT IS " FOCUSING " OR " DEFOCUSING "

- ② PROVE THAT THE " DEFOCUSING " NLS EQUATION $iu_t + u_{xx} - |u|^2 u = 0$ HAS NO SOLITARY WAVE SOLUTION (i.e. $u(x, t) = \exp[i(ax + bt)]f(x - ct)$) IN $L_2(\mathcal{R})$ W.R.T. x .
- ③ FIND TWO CONSERVATION LAWS OF THE 3WRI EQUATION
- ④ SOLVE THE STATIONARY 3WRI EQUATION IN TERMS OF ELLIPTIC FUNCTIONS